


$$\mathcal{L} = -mc^2 \sqrt{1-\beta^2} - V$$

START WITH NR LAGRANGIAN

$$\mathcal{L} = \underbrace{\frac{1}{2} m \dot{\vec{x}}^2}_{\text{NOT LI}} + \underbrace{q \dot{\vec{x}} \cdot \vec{A} - q\phi}_{\text{LORENTZ INV}}$$

NOT LI

LORENTZ INV

TOTAL DYNAMICAL ENERGY γmc^2

ELEM USE $\frac{d}{dt}$, WANT $\frac{d}{d\tau}$

$\Rightarrow$ DIVIDE BY $\gamma = (1 - \dot{\vec{x}}^2/c^2)^{-1/2}$

SUBTRACT THE REST ENERGY

$$\frac{\gamma mc^2 - mc^2}{\gamma} = mc^2 - mc^2 (1 - \beta^2)^{1/2}$$

NEGLECT CONSTANT TERM mc^2

$$\Rightarrow -mc^2 (1 - \beta^2)^{1/2}$$

COVARIANT FORMULATION

IDEALLY τ IS AN INDEPENDENT COORDINATE PARAMETER

$U \cdot U = U^\mu U_\mu = c^2$ THE SET OF $\{U^\nu\}$
ARE NOT INDEPENDENT

IGNORE THIS CONSTRAINT AND
APPLY IT LATER

$$x'^\mu = \frac{dx^\mu}{d\tau} \quad \dot{x}^\mu = \frac{dx^\mu}{dt}$$

$$\delta I = \delta \int_{\tau_1}^{\tau_2} \Lambda(x^\mu, x'^\mu) d\tau$$

WANT Λ THAT GIVES US THE EOM

$$I = \int_{t_1}^{t_2} \mathcal{L}(x^j, \dot{x}^j) dt = \frac{1}{c} \int_{\tau_1}^{\tau_2} \mathcal{L}(x^j, c \frac{x'^j}{x'^0}) x'^0 d\tau$$

$$\Lambda(x^\mu, x'^\mu) = \frac{x'^0}{c} \mathcal{L}(x^\mu, c \frac{x'^j}{x'^0})$$

$$\Lambda(x^\mu, a x'^\mu) = a \Lambda(x^\mu, x'^\mu)$$

$$\Lambda = x'^\mu \frac{\partial \Lambda}{\partial x'^\mu}$$

$$\frac{d}{d\tau} \left(\frac{\partial \Lambda}{\partial U^\nu} \right) - \frac{\partial \Lambda}{\partial x^\nu} = 0 \quad U^\nu = \frac{dx^\nu}{d\tau}$$

ONE POSSIBLE FREE PARTICLE Λ

$$\Lambda = \frac{1}{2} m U_\nu U^\nu$$

$$\Lambda = \frac{1}{2} m U_\nu U^\nu + g U^\mu A_\mu(x^\mu)$$

$$\frac{d}{dt}(mu^\nu) = -q \frac{dA^\nu}{dt} + \frac{\partial}{\partial x_\nu} (g_{\nu\mu} A_\mu)$$

THIS IS WHAT WE GOT BEFORE

$$\frac{dp^\mu}{d\tau} = K^\mu \quad \text{MINKOWSKI FORCE}$$

$$K_\mu = -q \left(\frac{\partial (u_\nu A^\nu)}{\partial x^\mu} - \frac{dA_\mu}{d\tau} \right)$$

$$F_i = K_i \sqrt{1-\beta^2}$$

MECHANICAL MOMENTUM $p^\mu = mu^\mu$

CANONICAL MOMENTUM

$$P^\mu = g^{\mu\nu} \frac{\partial \mathcal{L}}{\partial u^\nu} = \frac{\partial \mathcal{L}}{\partial u^\nu} = mu^\mu + qA^\mu$$

$$= p^\mu + qA^\mu$$

$$P^0 = \frac{E}{c} + q \frac{\phi}{c} = \frac{1}{c} \bar{E} \quad \begin{matrix} \text{MECHANICAL} \\ \text{ENERGY} \end{matrix} \quad \begin{matrix} \text{TOTAL} \\ \text{ENERGY} \end{matrix}$$

$$\vec{P} = \vec{p} + q\vec{A}$$

$$E^2 = c^2 (\vec{P} - q\vec{A})^2 + m^2 c^4$$

DOES NOT WORK FOR INTERNAL
~~FORCES~~ POTENTIALS

ACTION AT A DISTANCE PROBLEMS
DUE TO LACK OF RETARDED POT.

RELATIVISTIC HAMILTONIANS

NON COVARIANT, WORKS IN A SPECIFIC
LORENTZ FRAME
TIME AS A PARAMETER, NOT A COORDINATE

$$\mathcal{L} = -mc^2 \sqrt{1-\beta^2} - V$$

$$H = T + V \quad \text{TOTAL ENERGY}$$

$$T^2 = p^2 c^2 + m^2 c^4 \quad (T \text{ IS NOW MASS ENERGY PLUS KE})$$

$$H = \sqrt{p^2 c^2 + m^2 c^4} + V$$

$$\text{EM LAGRANGIAN } \mathcal{L} = -mc^2 \sqrt{1-\beta^2} + q \vec{A} \cdot \vec{v} - q\phi$$

$$\Rightarrow H = T + q\phi$$

$$\vec{P}^i = m \vec{v}^i + q \vec{A}^i \quad \text{CANONICAL MOMENTUM}$$

$$T^2 = (\vec{P} - q \vec{A})^2 c^2 + m^2 c^4$$

$$H = \left[(\vec{P} - q \vec{A})^2 c^2 + m^2 c^4 \right]^{1/2} + q\phi$$

$$\frac{H - q\phi}{c} = P^0 \quad \text{zeroth component of } m \vec{v}^\nu + q \vec{A}^\nu$$

H IS NOT INVARIANT. IT IS NOT A
LORENTZ SCALAR. IT TRANSFORMS
LIKE THE 0TH COMP. OF A 4-VECTOR

$$\text{WE STILL HAVE } \frac{\partial H}{\partial q_j} = \dot{p}_j \quad \frac{\partial H}{\partial p_j} = \dot{q}_j$$

COVARIANT APPROACH

x^μ IS A COORDINATE. IT HAS A CONJUGATE MOMENTUM

ASSUME $\exists \theta(\tau)$ WHICH IS MONOTONIC IN τ

$$\dot{q}_i = \frac{\partial q_i}{\partial \theta}$$

$$\Rightarrow \Lambda(q, q', t, t') = t' \mathcal{L}(q, \frac{q'}{t'}, t)$$

q, q' ARE $\{\dot{q}_i\} \{\dot{q}_i'\}$

CANONICAL MOMENTUM CONJUGATE TO t
(USE $\dot{q}_i = \frac{q_i'}{t'} = \frac{\partial q_i}{\partial \theta} / \frac{\partial t}{\partial \theta}$)

$$P_t = \frac{\partial \Lambda}{\partial t'} = \mathcal{L} + t' \frac{\partial \mathcal{L}}{\partial t'}$$

$$= \mathcal{L} - \frac{q_i'}{t'} \frac{\partial \mathcal{L}}{\partial \dot{q}_i} = \mathcal{L} - \dot{q}_i \frac{\partial \mathcal{L}}{\partial \dot{q}_i} = -H$$

$$P_i = \frac{\partial \Lambda}{\partial q_i'} = t' \frac{\partial \mathcal{L}}{\partial q_i'} = t' \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \frac{1}{t'} \right) = P_i \checkmark$$

OTHER MOMENTA UNCHANGED

BUT THE COVARIANT $\mathcal{L}$ WE USED IS
1st ORDER HOMOGENOUS IN q'

$$\Rightarrow H = 0$$

$$\text{TRY } \Lambda(x^\mu, u^\mu) = \frac{1}{2} m u_\mu u^\mu \quad (\text{FREE})$$

THIS GIVES CORRECT FORM

$$H = \dot{q} P - \mathcal{L} = P_\mu P^\mu / 2m$$

$$E+m: \quad \Lambda(x^\mu, u^\mu) = \frac{1}{2} m u_\mu u^\mu + g u^\mu A_\mu(x^\mu)$$

$$P_\mu = m u_\mu + g A_\mu$$

$$H = \frac{(P_\mu - g A_\mu)(P^\mu - g A^\mu)}{2m}$$

8 EQ OF MOTION
6 INDEPENDENT

$$\frac{dx^\mu}{d\tau} = \frac{\partial H}{\partial p^\mu} g^{\mu\nu} \quad \frac{dp^\mu}{d\tau} = - \frac{\partial H}{\partial x^\mu} g^{\mu\nu}$$

$$u^0 = \frac{\partial H}{\partial p^0} = \frac{1}{m} (p^0 - g A^0)$$

$$p^0 = \frac{1}{c} (T + g \phi) = \frac{H}{c}$$

$$\frac{dp^0}{d\tau} = \frac{1}{\sqrt{1-\beta^2}} \frac{dp^0}{dt} = - \frac{1}{c} \frac{\partial H}{\partial t}$$

$$\frac{dH}{dt} = \sqrt{1-\beta^2} \frac{\partial H}{\partial t}$$

REALLY TOUGH TO DO WITH OTHER
POTENTIALS

EVEN TOUGHER WITH INTERNAL
(PARTICLE - PARTICLE) POTENTIALS